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*Acta Cryst.* (1999). **C55**, 1753–1755

## Tl<sub>2</sub>MoO<sub>4</sub> at 350 K

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(Received 31 March 1999; accepted 25 June 1999)

### Abstract

The structure of the title compound, dithallium molybdate, is isotypical with the glaserite structure and Tl<sub>2</sub>WO<sub>4</sub>. The Mo ion is coordinated tetrahedrally by oxygen, with bond lengths ranging from 1.699(16) to 1.725(16) Å, and angles between 107.6(4) and 111.3(4)°. The shortest O–O edge length is 2.83(2) Å. Three independent Tl ions are observed in the structure. The coordination polyhedron around Tl1 can be described as a distorted octahedron, with Tl–O distances of 2.769(10) Å. Tl2 has a coordination number of 12, with six O atoms at a distance of 2.941(10) Å and the remaining six at a distance of 3.622(10) Å. The polyhedron around Tl3 is formed by ten O atoms, with distances ranging between 2.467(16) and 3.682(16) Å. The observed bond distances in Tl<sub>2</sub>MoO<sub>4</sub> are in good agreement with those in Tl<sub>2</sub>WO<sub>4</sub>.

### Comment

Tl<sub>2</sub>MoO<sub>4</sub> was first investigated by Gaultier & Pannetier (1972), who found three phase transitions, presumably all of first order, at temperatures of 776, 673 and 311 K. The lattice parameters given at room temperature are  $a = 7.919(5)$ ,  $b = 11.026(7)$  and  $c = 6.179(5)$  Å, in space group *Pnam*. Above 323 K, the structure is supposed to be trigonal in space group *P3̄m*, with lattice parameters  $a = 6.257(3)$  and  $c = 8.085(4)$  Å. These authors further reported that the phase transition at 673 K leads to the destruction of the crystals. Later investigations (Sleight *et al.*, 1975) reaffirmed that Tl<sub>2</sub>MoO<sub>4</sub> at 298 K is orthorhombic, although the authors reported the non-centrosymmetric space group *Pna2<sub>1</sub>* with similar lattice parameters. They could only observe the phase transitions at 310 and 772 K and did not detect that at 673 K. They also presumed the intermediate phase to have space group *P3̄m* and proposed the space group *P6<sub>3</sub>/mmc* for the high-temperature phase.

Our investigations corroborate the space group *P3̄m* for the intermediate phase, which we measured at 350 K. The refinement of the structure shows it to be isostructural with glaserite, K<sub>3</sub>Na(SO<sub>4</sub>)<sub>2</sub> (Okada & Oosaka, 1980), and with Tl<sub>2</sub>WO<sub>4</sub> (Okada *et al.*, 1979).

In Tl<sub>2</sub>MoO<sub>4</sub> (Fig. 1), the Mo ions are coordinated tetrahedrally by the O atoms, with bond lengths ranging from 1.699(16) to 1.725(16) Å, and angles between 107.6(4) and 111.3(4)°.

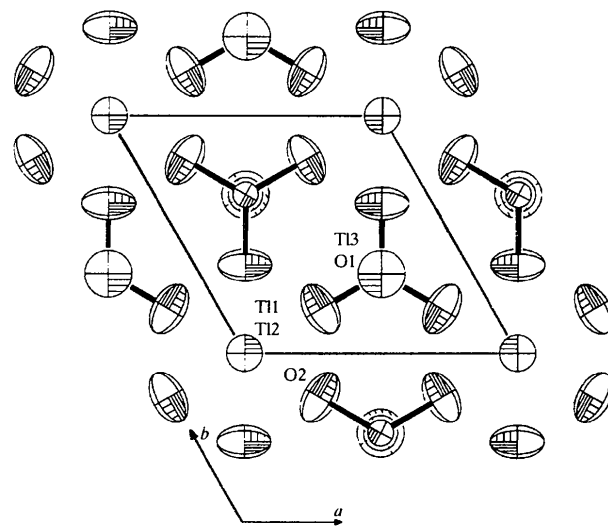


Fig. 1. The *a, b* projection of the structure of Tl<sub>2</sub>MoO<sub>4</sub>. Displacement ellipsoids are drawn at the 50% probability level.

As pointed out by Moore (1973), in the glaserite structure, and consequently in Tl<sub>2</sub>MoO<sub>4</sub>, three different monovalent cations can be observed. These are now discussed.

The coordination around Tl1 can be described as a distorted octahedron, with a Tl—O distance of 2.769(10) Å. This compares favourably with the distances observed in Tl<sub>2</sub>WO<sub>4</sub> [2.70(7) Å].

Tl2 has a coordination number of 12. Six O atoms are at a distance of 2.941(10) Å, while the remaining six are at a distance of 3.622(10) Å. For Tl<sub>2</sub>WO<sub>4</sub>, the corresponding distances are 2.95(7) and 3.630(4) Å. The coordination polyhedron can be described as a 12-coordinate trigonal antiprism, with six additional meridional anions in a hexagonal outline. The ideal point symmetry is  $\bar{3}m$ .

The polyhedron around Tl3 is formed by ten O atoms. It has ideal point symmetry of  $3m$ . Again, six of them (O2) form a distorted meridional hexagon, while the three O atoms above share a tetrahedral base with, and the one O atom below forms the apex of, MoO<sub>4</sub> tetrahedra. The Tl3—O1 distance to this apical oxygen is the shortest, at 2.467(16) Å, followed by the six Tl3—O2 distances of the hexagon [3.178(16) Å]. The largest distances are formed with the basal O atoms [3.682(16) Å]. These distances compare well with those observed in Tl<sub>2</sub>WO<sub>4</sub>: the shortest Tl3—O distance in this compound is 2.46(6) Å, followed by six distances of 3.18(6) Å and three more of 3.64(6) Å.

The shortest O—O distance is found in the Tl3 coordination polyhedron, with a value of 2.78(2) Å (Tl<sub>2</sub>WO<sub>4</sub> 2.87 Å), while the shortest O—O edge length in the tetrahedron is 2.83(2) Å (Tl<sub>2</sub>WO<sub>4</sub> 2.92 Å).

Bond valence sums, which were calculated according to Brese and O'Keefe (1991), show a too high value for Tl1 in the two compounds (Tl<sub>2</sub>MoO<sub>4</sub> 1.20; Tl<sub>2</sub>WO<sub>4</sub> 1.44). For the other two Tl ions, the values are slightly too low in the molybdate (Tl2 0.87; Tl3 0.90) and nearly ideal in the tungstate (Tl2 0.96; Tl3 0.97).

## Experimental

Crystals of the title compound were grown from an aqueous solution of Tl<sub>2</sub>CO<sub>3</sub> and MoO<sub>3</sub>, kept at 313 K for about one month. As the resulting crystals were very thin hexagonal plates, they were placed in a Pt crucible, heated to 923 K and afterwards cooled slowly to 873 K at a rate of 1 K d<sup>-1</sup>. After rapid cooling to room temperature, the crystals were of a suitable size for single-crystal investigations.

### Crystal data

Tl <sub>2</sub> MoO <sub>4</sub>	Mo K $\alpha$ radiation
$M_r = 568.68$	$\lambda = 0.71073$ Å
Trigonal	Cell parameters from 2313 reflections
$P\bar{3}m1$	$\theta = 2-26^\circ$
$a = 6.266(1)$ Å	$\mu = 60.489$ mm <sup>-1</sup>
$c = 8.103(2)$ Å	$T = 350(1)$ K
$V = 275.52(9)$ Å <sup>3</sup>	Irregular
$Z = 2$	$0.10 \times 0.10 \times 0.05$ mm
$D_x = 6.855$ Mg m <sup>-3</sup>	Yellow
$D_m$ not measured	

### Data collection

Stoe IPDS diffractometer	$R(F^2)_{\text{int}} = 0.115$
Imaging plate scans	$\theta_{\text{max}} = 26^\circ$
Absorption correction:	$h = -7 \rightarrow 7$
Gaussian (X-RED; Stoe, 1998)	$k = -7 \rightarrow 7$
$T_{\text{min}} = 0.005$ , $T_{\text{max}} = 0.041$	$l = -9 \rightarrow 9$
2008 measured reflections	50 standard reflections
237 independent reflections	intensity decay: 1%

### Refinement

Refinement on $F^2$	$\Delta\rho_{\text{max}} = 1.31$ e Å <sup>-3</sup>
$R[F^2 > 2\sigma(F^2)] = 0.039$	$\Delta\rho_{\text{min}} = -1.08$ e Å <sup>-3</sup>
$wR(F^2) = 0.092$	Extinction correction:
$S = 1.053$	SHELXL93 (Sheldrick, 1993)
237 reflections	Extinction coefficient:
21 parameters	0.0031(10)
$w = 1/[\sigma^2(F_o^2) + (0.0468P)^2]$	Scattering factors from
where $P = (F_o^2 + 2F_c^2)/3$	International Tables for
$(\Delta/\sigma)_{\text{max}} < 0.001$	Crystallography (Vol. C)

Table 1. Fractional atomic coordinates and equivalent isotropic displacement parameters (Å<sup>2</sup>)

$$U_{\text{eq}} = (1/3)\sum_i \sum_j U^{ij} a^i a^j a_i a_j$$

	x	y	z	$U_{\text{eq}}$
Tl1	0	0	0	0.0444(5)
Tl2	0	0	1/2	0.0620(6)
Tl3	2/3	1/3	0.17374(10)	0.0568(5)
Mo	2/3	1/3	0.6878(2)	0.0346(6)
O1	2/3	1/3	0.478(2)	0.094(7)
O2	0.1852(9)	-0.1852(9)	0.2350(15)	0.100(4)

Table 2. Selected geometric parameters (Å, °)

Tl1—O2 (× 6)	2.769(10)	Tl3—O2' (× 3)	3.682(16)
Tl2—O2 (× 6)	2.941(10)	Mo—O1	1.699(16)
Tl2—O1 (× 6)	3.622(10)	Mo—O2 <sup>ii</sup> (× 3)	1.725(16)
Tl3—O1	2.467(16)	O1—O2 <sup>iii</sup>	2.83(2)
Tl3—O2 (× 6)	3.178(16)	O2—O2 <sup>iiii</sup>	2.78(2)
O2—Tl1—O2 <sup>i</sup>	180.0(4)	O2—Tl2—O1 <sup>viii</sup>	130.3(4)
O2—Tl1—O2 <sup>i'</sup>	77.9(4)	O2—Tl2—O1 <sup>ix</sup>	112.1(4)
O2—Tl1—O2 <sup>ii</sup>	102.1(4)	O2—Tl2—O1 <sup>x</sup>	49.7(4)
O2—Tl2—O2 <sup>i</sup>	72.6(4)	O1—Mo—O2 <sup>ii</sup>	111.3(4)
O2—Tl2—O2 <sup>iii</sup>	107.4(4)	O2 <sup>ii</sup> —Mo—O2 <sup>ii</sup>	107.6(4)
O2—Tl2—O1	67.9(4)		

Symmetry codes: (i)  $1-x, -y, -z$ ; (ii)  $1-x, -y, 1-z$ ; (iii)  $1-x+y, -x, z$ ; (iv)  $-x, -y, -z$ ; (v)  $-y, x-y, z$ ; (vi)  $y, -x+y, -z$ ; (vii)  $y, -x+y, 1-z$ ; (viii)  $x-1, y, z$ ; (ix)  $1-x, 1-y, 1-z$ ; (x)  $x-y, x, 1-z$ .

The computations were carried out with the coordinates of Tl<sub>2</sub>WO<sub>4</sub> (Okada *et al.*, 1979) as the starting model. Due to the high anisotropic displacement parameters of the O atoms, a refinement with split positions for these atoms was attempted. However, this did not lead to a significantly better model and so this refinement was discarded.

Data collection: EXPOSE (Stoe, 1997). Cell refinement: EXPOSE. Data reduction: X-RED (Stoe & Cie, 1998). Program(s) used to refine structure: SHELXL93 (Sheldrick, 1993). Molecular graphics: ORTEP-3 for Windows (Farrugia, 1996).

The authors, especially KF, gratefully acknowledge financial support from the European Union (TMR Project ERBFMBICT96-1527), the Gobierno Vasco (PI97/71) and the UPV (EB098/97).

Supplementary data for this paper are available from the IUCr electronic archives (Reference: BR1242). Services for accessing these data are described at the back of the journal.

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*Acta Cryst.* (1999). **C55**, 1755–1757

## Orthorhombic $\text{InFe}_{0.33}\text{Ti}_{0.67}\text{O}_{3.33}$

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(Received 6 April 1999; accepted 1 July 1999)

## Abstract

The title compound, indium iron titanium oxide, is closely related to  $\text{InFeO}_3$ , having a hexagonal structure which consists of alternating layers of  $\text{InO}_6$  octahedra and  $\text{FeO}_5$  trigonal bipyramids. According to substitution of  $\text{Ti}^{4+}$  for  $\text{Fe}^{3+}$ , excess O atoms are introduced into the Fe–O trigonal lattice plane of  $\text{InFeO}_3$ . The in-plane arrangement of O atoms can be described as partial occupation on a honeycomb lattice, although large displacement parameters indicate local shifts of O atoms due to repulsive interactions between them.

## Comment

Recently, new phases with monoclinic and orthorhombic structures have been found in the system  $\text{InFeO}_3\text{--In}_2\text{Ti}_2\text{O}_7$  (Brown, Flores *et al.*, 1999), and their isostruc-

tural compounds were also obtained in various solid-solution systems (Brown, Kimizuka *et al.*, 1999). Structural details of these phases are of interest but complicated due to commensurate or incommensurate satellite reflections (Brown, Flores *et al.*, 1999). This study is the first on the analysis of this family. The average structure of orthorhombic  $\text{InFe}_{0.33}\text{Ti}_{0.67}\text{O}_{3.33}$  has been determined using main reflections only, although the crystals showed incommensurate ( $q \approx 0.31a^*$ ) satellite reflections.

The structure is closely related to that of  $\text{InFeO}_3$ , in which two kinds of coordination polyhedral layers are stacked along the *c* axis; one is an  $\text{InO}_6$  octahedral layer and the other is an  $\text{FeO}_5$  trigonobipyramidal layer (Giaquinta *et al.*, 1994). The structure of orthorhombic  $\text{InFe}_{0.33}\text{Ti}_{0.67}\text{O}_{3.33}$  is given by deformation of the hexagonal  $\text{InFeO}_3$  structure [ $a = 3.3270(2) \text{ \AA}$  and  $c = 12.1750(1) \text{ \AA}$ ]. A considerable amount of  $\text{Ti}^{4+}$  is substituted for  $\text{Fe}^{3+}$  in  $\text{InFe}_{0.33}\text{Ti}_{0.67}\text{O}_{3.33}$ . The charge neutrality is maintained by the introduction of excess O atoms into the Fe–O2 trigonal plane. This phase being clearly distinct from  $\text{InFeO}_3$ , the O3 site, as well as the O2 site, is used to accommodate O atoms. Occupation ratios are greater than 0.5 at both sites and the amount of O atoms in the plane is 1.33 times that of  $\text{InFeO}_3$ . Thus, the arrangement of O atoms in this plane results in a partially occupied honeycomb lattice, as illustrated in Fig. 1(a), while  $\text{InFeO}_3$  gives a fully occupied triangle lattice of O2 sites, as shown in Fig. 1(b). An example of prob-

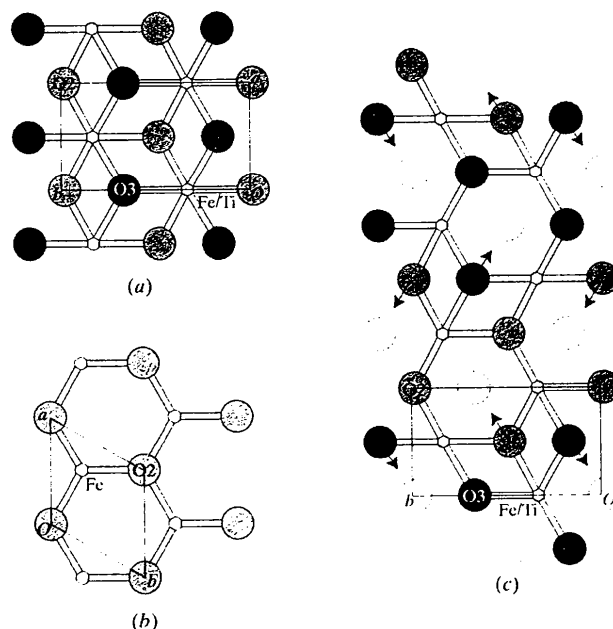


Fig. 1. The atomic arrangement on the plane  $z = \frac{1}{4}$  for (a) the average structure of  $\text{InFe}_{0.33}\text{Ti}_{0.67}\text{O}_{3.33}$  with the O2 and O3 sites partially occupied, (b)  $\text{InFeO}_3$  and (c) an example of local structures of  $\text{InFe}_{0.33}\text{Ti}_{0.67}\text{O}_{3.33}$ . Dotted circles are vacant oxygen sites. Arrows indicate probable displacements of O atoms.